



Vibrational spectroscopy and microwave dielectric properties of two novel $\text{Ca}_3\text{Ln}_2\text{W}_2\text{O}_{12}$ (Ln = La, Sm) tungstate ceramics

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ABSTRACT

Two novel tungstate ceramics with a nominal composition of $\text{Ca}_3\text{Ln}_2\text{W}_2\text{O}_{12}$ (Ln = La, Sm) were prepared via standard solid-state reaction methods. According to X-ray diffraction patterns, both ceramics crystallized in hexagonal crystal systems with a space group of $R\bar{3}m$ (No. 166). The vibrational modes of the Raman spectra were identified, and the evolutions of the wavenumber and the full width at half maximum (FWHM) were analyzed. The optimum microwave dielectric properties with $\epsilon_r = 18.7$, $Qf = 50,500$ GHz, and $\tau_f = -90$ ppm/°C were realized in $\text{Ca}_3\text{La}_2\text{W}_2\text{O}_{12}$ ceramics, while moderate Qf values were obtained in the Sm composition ($\epsilon_r = 19.5$, $Qf = 15,700$ GHz, and $\tau_f = -95$ ppm/°C). The discrepancy in the Qf values was associated with the broader FWHM of Raman spectra and smaller grain sizes. Infrared reflectivity (IR) spectra were also used to determine the intrinsic dielectric loss, where more damped phonon parameters were also identified in $\text{Ca}_3\text{Sm}_2\text{W}_2\text{O}_{12}$ ceramics.

1. Introduction

Over the past decades, in pace with the booming industry of microwave telecommunications, microwave dielectric ceramics have become indispensable components of various microwave devices due to their compactness, light weight, thermal stability, low cost and excellent performance [1]. From the device design perspective, dielectrics with suitable values of the dielectric constant (ϵ_r) are required for balancing device miniaturization and short signal delay time. A high Qf value ($Q = 1/\tan\delta$ and f denotes the resonant frequency) can reduce the energy loss during signal propagation and create satisfactory frequency selectivity, which is of vital importance to minimizing the signal attenuation. In addition, a low temperature coefficient of resonant frequency (τ_f) is required for ensuring frequency stability. Last, for massive industrial production, the cost effectiveness is another important factor as the typical ultra-low-loss candidates, such as Ba-based complex perovskites, typically contain noble raw materials, such as tantalates or niobates [2,3]. The identification of a single material that satisfies all these requirements is a formidable task, and optimally balancing these properties is a major challenge in the microwave dielectric ceramics industry.

Tungstates have been regarded as important candidates for dielectric materials due to their low cost, facile synthesis, satisfactory chemical stability, and excellent dielectric properties. Among tungstates, AWO_4 (A = Ca, Sr, Ba, Zn, Co, or Ni) compounds with

monoclinic wolframite or tetragonal scheelite structures are well known [4–7]. Pullar et al. reported that AWO_4 ceramics that are sintered at 1200 °C typically exhibit low dielectric constant values of approximately 12 and Qf values in the range of 24,900–62,800 GHz [4]. In 2011, Li_2WO_4 ceramics with an ultra-low sintering temperature of approximately 650 °C were reported by Zhou et al., and their satisfactory microwave dielectric properties ($\epsilon_r \sim 5.5$, $Qf \sim 62,000$ GHz, and $\tau_f = -146$ ppm/°C) indicated high potential in ultra-low-temperature cofired ceramic (ULTCC) technologies [8]. More recently, various types of tungstates such as Li_4WO_5 , LiAlW_2O_8 , and $\text{Li}_2\text{Mg}_2(\text{WO}_4)_3$ were proposed, and their potential applications in microwave dielectric ceramics were systematically investigated [9–11]. $\text{Ca}_3\text{Ln}_2\text{W}_2\text{O}_{12}$ (Ln = La and Sm) powders are widely applied as promising host materials for inorganic phosphors due to their outstanding structural tunability and luminescent properties [12–14]. However, to the best of our knowledge, minimal investigation has been conducted on the potential applications of $\text{Ca}_3\text{Ln}_2\text{W}_2\text{O}_{12}$ ceramics or on their microwave dielectric properties. Therefore, the exploration of the structural parameters and microwave dielectric properties of $\text{Ca}_3\text{Ln}_2\text{W}_2\text{O}_{12}$ ceramics could be important, and may trigger a new breakthrough in the development of novel tungstate-based microwave dielectric ceramics that realize high performances.

In this study, the crystal structures and Raman spectra of $\text{Ca}_3\text{Ln}_2\text{W}_2\text{O}_{12}$ ceramics are systematically investigated, together with their effects on the microwave dielectric properties. Moreover, IR

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spectra are extrapolated to the microwave frequency region to further identify the intrinsic dielectric properties.

2. Experimental procedure

$\text{Ca}_3\text{Ln}_2\text{W}_2\text{O}_{12}$ ceramics were prepared via high temperature solid-state reactions of CaCO_3 (99.99%), La_2O_3 (99.99%), Sm_2O_3 (99.99%) and WO_3 (99.99%) raw powders. Prior to weighing, La_2O_3 and Sm_2O_3 were preheated at 900 °C for 2 h to remove the moisture. The stoichiometric powder was ball-milled with zirconia media at a speed of 180 r/min for 6 h. After that, the obtained mixtures were calcined at 1100 °C for 3 h and pressed into cylindrical pellets of 12 mm in diameter and 5 mm in height. Finally, the pellets were sintered at 1300 °C–1400 °C for 3 h at a heating rate of 5 °C/min to obtain the dense ceramics.

X-ray diffraction patterns were collected using a RIGAKU D/max 2550/PC. Scanning electron micrographs were obtained from the polished and thermally etched surfaces using a SIRION-100 system. Raman spectra were recorded using an HR-800 LabRaman device. IR spectra were collected using an IFS 66v/s infrared spectrometer. The relative densities were measured via the Archimedes method. The microwave dielectric properties were evaluated using a silver-coated resonant cavity connected to a vector network analyzer (E8363B, Agilent Technologies Inc., Palo Alto, CA) [15,16]. The measurement of τ_f value was conducted in the temperature range of 20–80 °C.

3. Results and discussion

The XRD patterns of $\text{Ca}_3\text{Ln}_2\text{W}_2\text{O}_{12}$ ceramics that were sintered at 1300 °C are presented in Fig. 1(a). Both patterns exhibit similar diffraction features, and all the diffraction peaks can be well indexed according to the standard PDF card of $\text{Ca}_3\text{La}_2\text{W}_2\text{O}_{12}$ (JCPDS #49-0965). Hence, monophasic $\text{Ca}_3\text{Ln}_2\text{W}_2\text{O}_{12}$ ceramics have been successfully obtained in this study. Based on the cell refinements, the cell parameters of $\text{Ca}_3\text{Ln}_2\text{W}_2\text{O}_{12}$ ceramics are calculated as follows: $a = b = 9.76375$ Å, $c = 55.54702$ Å, $\gamma = 120^\circ$ for La and $a = b = 9.75646$ Å, $c = 55.47817$ Å, and $\gamma = 120^\circ$ for Sm. Moreover, according to the enlarged figure in the inset, the corresponding diffraction peaks of the Sm composition indicate higher 2θ than that of the La composition, which corresponds to cell shrinkage of the $\text{Ca}_3\text{Sm}_2\text{W}_2\text{O}_{12}$ ceramics. Although the crystal structure of $\text{Ca}_3\text{Ln}_2\text{W}_2\text{O}_{12}$ has not been reported yet, it is widely believed that it crystallizes in a hexagonal crystal system with a space group of $R\bar{3}m$, which is similar to that of $\text{Ca}_5\text{Re}_2\text{O}_{12}$ [17,18]. As approximately depicted in Fig. 1(b), the A-site cations (Ca^{2+} and Ln^{3+}) have four types of sites [coordination number (CN) = 6, 8, and 9], while three types of W^{6+} sites are observed and they are all located in the centre of the oxygen octahedra (CN = 6).

SEM images of the polished and thermally etched surfaces of $\text{Ca}_3\text{Ln}_2\text{W}_2\text{O}_{12}$ ceramics that were sintered at 1375 °C are shown in Fig. 2. Both samples exhibit dense microstructures with normal grain size distributions, and the grain exhibit a hexangular shape-like morphology. The average grain size of $\text{Ca}_3\text{La}_2\text{W}_2\text{O}_{12}$ (9.85 μm) is slightly larger than that of the Sm composition (8.26 μm). Moreover, no abnormal grain growth with a strong preferred orientation is observed, even though the crystal structure of the $\text{Ca}_3\text{Ln}_2\text{W}_2\text{O}_{12}$ ceramics corresponds to a large c/a value.

Room-temperature Raman spectra of the $\text{Ca}_3\text{Ln}_2\text{W}_2\text{O}_{12}$ ceramics are shown in Fig. 3(a). Strong similarities are observed when comparing the Raman spectra of the present ceramics with those of $\text{Ca}_3\text{La}_2\text{Te}_2\text{O}_{12}$, LiH_5TeO_6 , and Bi_2WO_6 , etc [18–20]. The Raman spectra of the $\text{Ca}_3\text{Ln}_2\text{W}_2\text{O}_{12}$ ceramics exhibit eight bands: a very strong high-frequency shoulder band at 780–860 cm^{-1} , two very weak and broad bands at 700–780 cm^{-1} and 480–540 cm^{-1} , two medium-intensity bands at approximately 445 and 287 cm^{-1} and three weak bands at approximately 368, 225, and 129 cm^{-1} . According to Poirier et al., the Raman bands at 500–900 cm^{-1} are assigned to the asymmetric and symmetric stretching vibrations of the bridging W–O–W bonds [20,21]. The two bands at approximately 445 and 368 cm^{-1} are typical characteristics of the bending vibrations of W–O terminal bonds in WO_6 octahedra [22]. The remaining three low-frequency bands may be associated with the nonsymmetric bending of the Ca/La–O bond [23]. With substitution of smaller Sm^{3+} ($r = 0.958$ Å, CN = 6; $r = 1.079$ Å, CN = 8; $r = 1.132$ Å, CN = 9) for La^{3+} ($r = 1.032$ Å, CN = 6; $r = 1.16$ Å, CN = 8; $r = 1.216$ Å, CN = 9) [24], the Raman spectrum of $\text{Ca}_3\text{Sm}_2\text{W}_2\text{O}_{12}$ undergoes an overall blueshift compared with that of the La composition. By simplifying the lattice vibration as a harmonic resonant model, the corresponding wavenumbers (ω) could be correlated with the bond strength (k) and the reduced mass (m^*) using the following equation [25]:

$$\omega = \sqrt{k/m^*} \quad (1)$$

Fig. 3(b) presents the wavenumber differences ($\Delta\omega = \omega_{\text{Sm}} - \omega_{\text{La}}$) of the labeled Raman peaks (peaks A–G). As the atomic mass of Sm is larger than that of La, the positive $\Delta\omega$ values in Fig. 3(b) indicate that the bond strength (k) of $\text{Ca}_3\text{Sm}_2\text{W}_2\text{O}_{12}$ ceramics should be stronger than that of the La composition. Therefore, it is reasonable to infer that the vibration bonds of $\text{Ca}_3\text{Ln}_2\text{W}_2\text{O}_{12}$ could be more compressed after the rational substitution of Sm for La. Moreover, FWHM of the Raman peak has been widely accepted as an important indicator of microwave dielectric properties, and the broadening of FWHM could be closely related to the decay of microwave propagation [26–28]. The FWHM differences of the present ceramics ($\Delta\text{FWHM} = \text{FWHM}_{(\text{Sm})} - \text{FWHM}_{(\text{La})}$) are presented in Fig. 3(c). Except for the split peak at ~ 820 cm^{-1} (peak G), all the Raman peaks of the $\text{Ca}_3\text{Sm}_2\text{W}_2\text{O}_{12}$ ceramics are

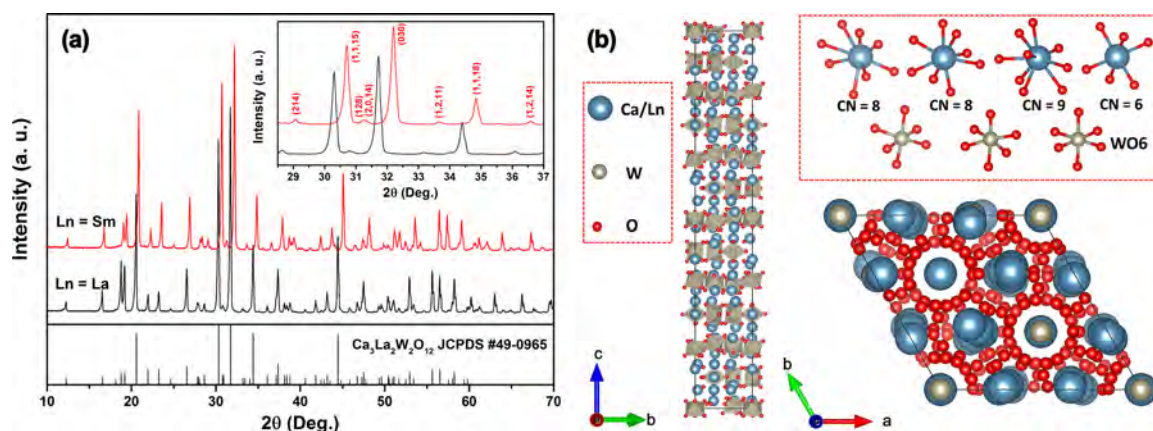


Fig. 1. (a) XRD patterns of $\text{CaLn}_2\text{W}_2\text{O}_{12}$ ceramics (Ln = La, Sm). (b) Schematic illustrations of the approximate crystal structure of $\text{Ca}_3\text{Ln}_2\text{W}_2\text{O}_{12}$ ceramics and the coordination environments of the cations.

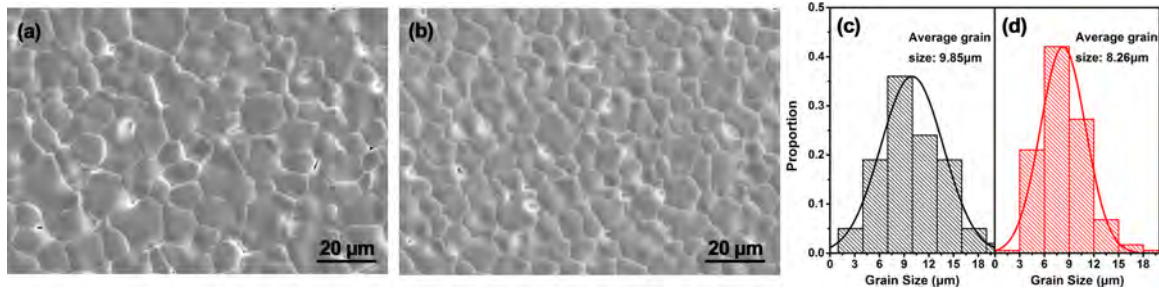


Fig. 2. SEM images and grain size distributions of (a,c) $\text{Ca}_3\text{La}_2\text{W}_2\text{O}_{12}$ and (b,d) $\text{Ca}_3\text{Sm}_2\text{W}_2\text{O}_{12}$ ceramics that were sintered at 1375 °C.

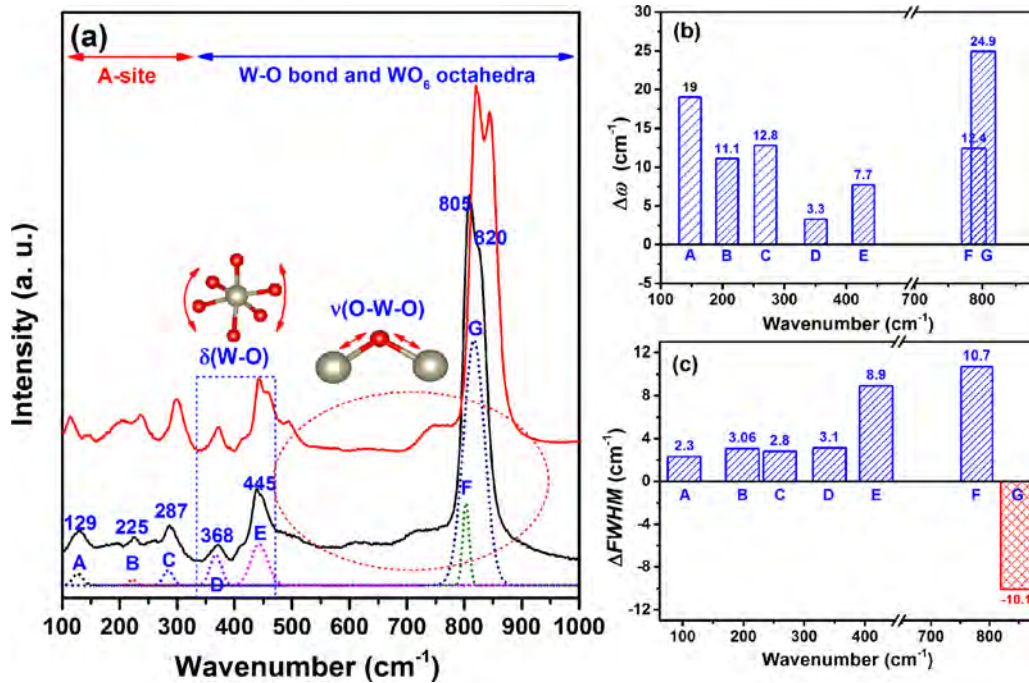


Fig. 3. (a) Raman spectra of $\text{Ca}_3\text{Ln}_2\text{W}_2\text{O}_{12}$ ceramics and the differences in the (b) wavenumber and (c) FWHM between La and Sm compositions. Note. Abbreviations: ν , stretching; δ , deformation, or bending.

widened compared with the corresponding peaks in the La composition.

The variations of the relative density and the microwave dielectric properties are presented in Fig. 4. The relative densities of both ceramics increase monotonously with increasing sintering temperature until they reach the maximum values of above 0.95 at 1375 °C and degrade slightly thereafter. The variation trends of the ϵ_r values are in satisfactory agreement with those of the relative densities. Hence, the porosity should be the dominant factor in controlling ϵ_r . To exclude the impact of porosity (P), the following empirical equation, which was proposed by Alford et al., is used to determine the actual dielectric constant (ϵ_m) [29]:

$$\epsilon_r = \epsilon_m \left(1 - \frac{3P(\epsilon_m - 1)}{2\epsilon_m + 1} \right) \quad (2)$$

The variation of the calculated ϵ_m values is plotted in Fig. 4(b). The ϵ_m values are located mainly at approximately 20 and vary indistinctively with increasing sintering temperature. This further demonstrates the important role of the relative density in determining the resultant dielectric constant. The Q_f values of $\text{Ca}_3\text{La}_2\text{W}_2\text{O}_{12}$ ceramics increase monotonously with the increasing sintering temperature until the maximum values of 50,500 GHz for $\text{Ca}_3\text{La}_2\text{W}_2\text{O}_{12}$ and 15,700 GHz for $\text{Ca}_3\text{Sm}_2\text{W}_2\text{O}_{12}$ are attained at 1375 °C. The τ_f values of the present ceramics vary around -90 ppm/°C and indicate slight downward variation trends with increasing sintering temperature (see Fig. 4(d)).

To further adjust the temperature stability, adding compounds with high positive τ_f value, such as TiO_2 (+450 ppm/°C) [30], should be effective and is now in progress. The optimal combinations of the microwave dielectric properties are identified as follows: $\epsilon_r = 18.7$, $Q_f = 50,500$ GHz, and $\tau_f = -90$ ppm/°C for $\text{Ca}_3\text{La}_2\text{W}_2\text{O}_{12}$ and $\epsilon_r = 19.5$, $Q_f = 15,700$ GHz, and $\tau_f = -95$ ppm/°C for $\text{Ca}_3\text{Sm}_2\text{W}_2\text{O}_{12}$.

Notoriously, the variation of the Q_f values can be explained by both intrinsic factors (such as structural symmetry and ionic/electronic polarizability) and extrinsic factors (microstructural flaws, such as porosity, grain boundary and secondary phase) [31,32]. For the extrinsic part, the effects of a secondary phase can be excluded as monophasic compositions have been demonstrated by the XRD results. Hence, the extrinsic dielectric loss should be mainly discussed in terms of the porosity, grain boundary, and electrical conductivity [33]. For each composition, the Q_f value indicates an overall improvement with increasing sintering temperature, which corresponds to the decline of the porosity. The larger average grain size of $\text{Ca}_3\text{La}_2\text{W}_2\text{O}_{12}$ ceramics should be a key extrinsic factor for their much higher Q_f value. The variation of electrical conductivity and its effect on the resultant Q_f values could be evaluated via measuring the impedance under various temperature, which needs to be further investigated. Furthermore, for the intrinsic part, the overall larger FWHM of $\text{Ca}_3\text{Sm}_2\text{W}_2\text{O}_{12}$ ceramics should correspond to a more damped lattice vibration, which could play an important role in their intrinsic dielectric loss.

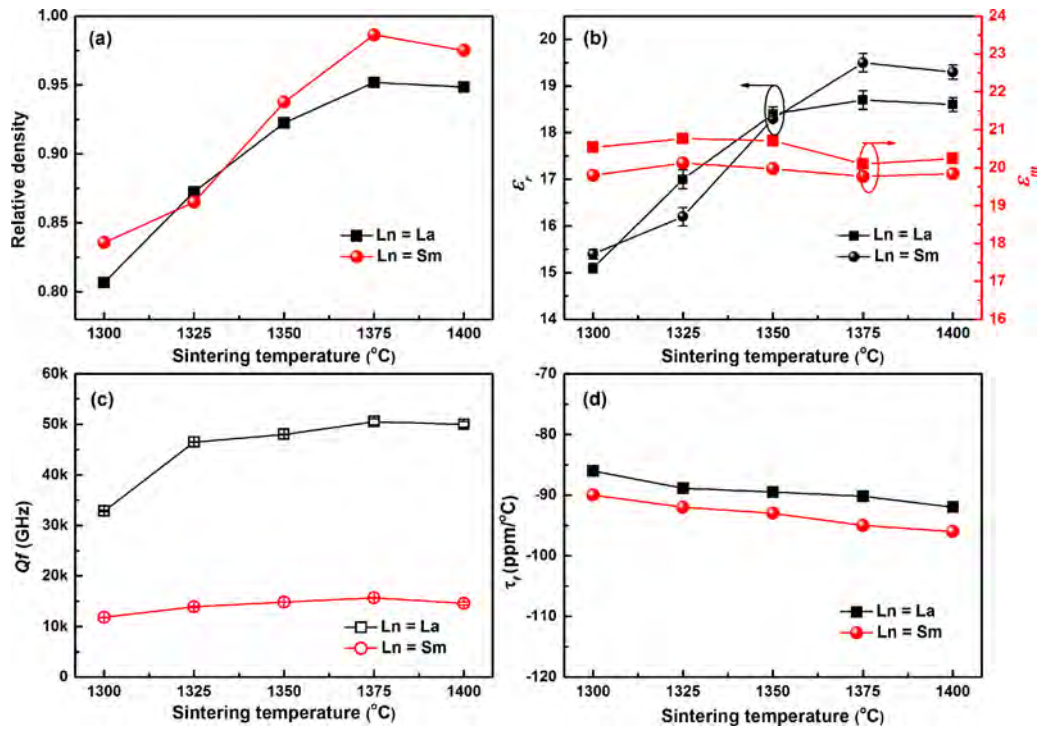


Fig. 4. (a) The relative density and (b-d) the microwave dielectric properties of $\text{Ca}_3\text{Ln}_2\text{W}_2\text{O}_{12}$ ceramics as functions of the sintering temperature.

Table 1

Phonon parameters that were calculated by fitting the IR spectra of $\text{Ca}_3\text{Ln}_2\text{W}_2\text{O}_{12}$ ($\text{Ln} = \text{La}, \text{Sm}$) ceramics.

La	$\omega_j(\text{cm}^{-1})$	$\gamma_j(\text{cm}^{-1})$	$\Delta\epsilon_j$	$\Delta\tan\delta_j \times 10^4$ at 10GHz	Sm	$\omega_j(\text{cm}^{-1})$	$\gamma_j(\text{cm}^{-1})$	$\Delta\epsilon_j$	$\Delta\tan\delta_j \times 10^4$ at 10GHz
1	177.0	42.6	5.2	1.1777	1	175.8	54.8	6.9	2.1982
2	210.0	58.1	2.2	0.3253	2	238.3	63.5	1.8	0.3246
3	303.0	39.7	2.5	0.1923	3	306.6	35.9	2.9	0.2078
4	330.3	31.0	2.4	0.1205	4	335.1	30.5	1.7	0.0812
5	363.4	33.5	0.7	0.0314	5	365.5	26.6	0.5	0.0119
6	468.9	46.6	0.3	0.0101	6	464.9	39.9	0.2	0.0077
7	573.4	38.7	0.6	0.0130	7	575.0	49.4	0.7	0.0193
8	628.5	120.6	0.5	0.0313	8	645.6	81.9	0.4	0.0090
9	820.2	101.8	0.01	0.0001	9	848.6	80.4	0.1	0.0013
ϵ_∞			4.2	0.0000	ϵ_∞			4.2	0.0000
Σ			18.61	1.9017	Σ			19.4	2.9303

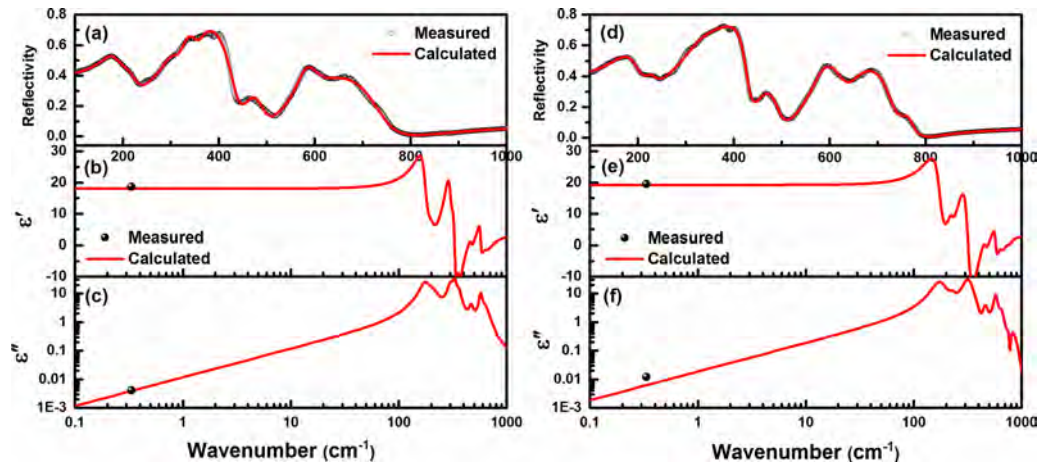


Fig. 5. Measured and calculated IR spectra of (a) $\text{Ca}_3\text{La}_2\text{W}_2\text{O}_{12}$ and (d) $\text{Ca}_3\text{Sm}_2\text{W}_2\text{O}_{12}$ ceramics and the real (ϵ') and imaginary parts (ϵ'') of the complex dielectric constants of (b,c) $\text{Ca}_3\text{La}_2\text{W}_2\text{O}_{12}$ and (e,f) $\text{Ca}_3\text{Sm}_2\text{W}_2\text{O}_{12}$ ceramics sintered at 1375 °C.

To further investigate the intrinsic dielectric properties, IR spectra of $\text{Ca}_3\text{Ln}_2\text{W}_2\text{O}_{12}$ ceramics sintered at 1375°C is fitted according to the phonon parameters that are listed in Table 1. The fitting method has been described in detail in our previous studies [25,31]. In the infrared frequency region, extrinsic factors will not affect the dielectric responses, and the calculated values from IR spectra can be used to distinguish the intrinsic dielectric properties from the measured dielectric properties. As shown in Fig. 5(a, d), the fitted spectra well accord with the measured spectra. Moreover, the calculated real (ϵ') and imaginary (ϵ'') parts of the complex dielectric constant are plotted and compared with the calculated values (black circles) (see Fig. 5(b,e) and (c,f)). The calculated and measured dielectric losses are of the same order of magnitudes (see Fig. 5(c,f)). Hence, the intrinsic dielectric loss should play a leading role in the overall dielectric losses. Moreover, the calculated intrinsic dielectric loss ($\Sigma \tan \delta_j$) of $\text{Ca}_3\text{Sm}_2\text{W}_2\text{O}_{12}$ (2.9303×10^{-4}) is much higher than that of the La composition (1.9017×10^{-4}), which corresponds to the lower measured Q_f values in $\text{Ca}_3\text{Sm}_2\text{W}_2\text{O}_{12}$ ceramics. According to Table 1, this is mainly due to the more damped phonon mode at $\sim 175\text{ cm}^{-1}$, and the contributions from higher frequency modes decrease rapidly with increasing ω_j since $\Delta \epsilon_j \sim S_j \omega_j^{-2}$ and $\Delta \tan \delta_j \sim \omega S_{fj} \omega_j^{-4}$ [34].

4. Conclusions

Two novel tungstate ceramics with a nominal composition of $\text{Ca}_3\text{Ln}_2\text{W}_2\text{O}_{12}$ (Ln = La, Sm) have been prepared via standard solid-state reaction methods. The XRD results demonstrate the formation of monophasic $\text{Ca}_3\text{Ln}_2\text{W}_2\text{O}_{12}$ ceramics with a space group of $R\bar{3}m$. Dense microstructures with a hexangular-shape-like grain morphology are obtained when sintered at 1375°C , where the optimal microwave dielectric properties ($\epsilon_r = 18.7$, $Q_f = 50,500\text{ GHz}$, and $\tau_f = -90\text{ ppm}/^\circ\text{C}$) are also realized in $\text{Ca}_3\text{La}_2\text{W}_2\text{O}_{12}$ ceramics. The ϵ_r and τ_f values of $\text{Ca}_3\text{Sm}_2\text{W}_2\text{O}_{12}$ ceramics are similar to those of the La composition, while their optimal Q_f value ($15,700\text{ GHz}$) is much lower. The discrepancy in the Q_f values is also demonstrated by the broader Raman peaks and more damped phonon parameters that are obtained by fitting the IR spectra.

Declaration of Competing Interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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